

BGG machinery and bounded Poincaré operators

Kaibo Hu

University of Oxford

based on joint papers with Andreas Čap (Vienna);
Snorre Christiansen (Oslo), Espen Sande (EPFL)

Structure-preserving numerical methods for partial differential
equations

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Three complexes

de Rham complex topology, curl, Poincaré inequality

$$0 \longrightarrow C^\infty \xrightarrow{\text{grad}} C^\infty \otimes \mathbb{R}^3 \xrightarrow{\text{curl}} C^\infty \otimes \mathbb{R}^3 \xrightarrow{\text{div}} C^\infty \longrightarrow 0.$$

elasticity complex Riemannian geometry, curl $\circ T \circ$ curl, Korn inequality

displacement formulation
/

Kröner's continuum description of dislocations/defects,
internal stress
┐

$$0 \longrightarrow C^\infty \otimes \mathbb{V} \xrightarrow{\text{def}} C^\infty \otimes \mathbb{S} \xrightarrow{\text{inc}} C^\infty \otimes \mathbb{S} \xrightarrow{\text{div}} C^\infty \otimes \mathbb{V} \longrightarrow 0$$

displacement

strain (metric)

stress (curvature)

force

\ intrinsic elasticity (Ciarlet et al.)

└ Hellinger-Reissner principle of elasticity

\mathbb{V} : vectors, \mathbb{S} : symmetric matrices

conformal complex conformal geometry, $\text{curl} \circ S^{-1} \circ \text{curl} \circ S^{-1} \circ \text{curl}$, conformal Korn inequality

ker of dev def: conformal Killing v.f. Cotton-York: flatness in conformal geometry

$$0 \longrightarrow H^s(\Omega) \otimes \mathbb{V} \xrightarrow{\text{dev def}} H^{s-1}(\Omega) \otimes (\mathbb{S} \cap \mathbb{T}) \xrightarrow{\text{Cott}} H^{s-4}(\Omega) \otimes (\mathbb{S} \cap \mathbb{T}) \xrightarrow{\text{div}} H^{s-5}(\Omega) \otimes \mathbb{V} \longrightarrow 0$$

gravitational wave variable: $\underbrace{\text{transverse-traceless (TT) gauge}}_{\text{(= symmetric, trace-free, div-free)}}$ $\underbrace{\text{stress like variable def } u \text{ in NS}}_{\text{(Gopalakrishnan, Lederer, Schöberl, 2019)}}$

\mathbb{T} : trace-free matrices

$$\text{dev } w := w - \frac{1}{n} \text{tr}(w)I, \quad \text{cott } g := \text{curl } S^{-1} \text{curl } S^{-1} \text{curl}, \quad \text{div } v := \nabla \cdot v$$

$$Su := u^T - \text{tr}(u)I$$

- NS: stress-like variable $\sigma := \text{def}(u)$ (Gopalakrishnan, Lederer, Schöberl 2020)
- transverse-traceless (TT) gauge of gravitational waves
- maximal slicing $K = 0$ (for constructing black hole initial data)
- gravity shielding, linearized Einstein constraint equations (Beig, Chruściel 2020)

Conformal complex: 2-rows approach (Arnold, KH 2021)

either eliminate trace from symmetric matrices:

$$\begin{array}{ccccccccccc}
 0 & \longrightarrow & H^q \otimes \mathbb{V} & \xrightarrow[\text{- mskw}]{\text{dev grad}} & H^{q-1} \otimes \mathbb{T} & \xrightarrow[\mathbb{S}]{\text{sym curl}} & H^{q-2} \otimes \mathbb{S} & \xrightarrow[\text{tr}]{\text{div div}} & H^{q-4} & \longrightarrow & 0 \\
 & & & \nearrow & & \nearrow & & \nearrow & & & \\
 0 & \longrightarrow & H^{q-1} \otimes \mathbb{V} & \xrightarrow{\text{def}} & H^{q-2} \otimes \mathbb{S} & \xrightarrow{\text{inc}} & H^{q-4} \otimes \mathbb{S} & \xrightarrow{\text{div}} & H^{q-5} \otimes \mathbb{V} & \longrightarrow & 0
 \end{array}$$

or eliminate skew-symmetric part from trace-free matrices:

$$\begin{array}{ccccccccccc}
 0 & \longrightarrow & H^{q-1} \otimes \mathbb{V} & \xrightarrow{\text{def}} & H^{q-2} \otimes \mathbb{S} & \xrightarrow{\text{inc}} & H^{q-4} \otimes \mathbb{S} & \xrightarrow{\text{div}} & H^{q-5} \otimes \mathbb{V} & \longrightarrow & 0 \\
 & & \nearrow \iota & & \nearrow \mathbb{S} & & \nearrow 2 \text{ vskw} & & & & \\
 0 & \longrightarrow & H^{q-2} & \xrightarrow{\text{hess}} & H^{q-4} \otimes \mathbb{S} & \xrightarrow{\text{curl}} & H^{q-5} \otimes \mathbb{T} & \xrightarrow{\text{div}} & H^{q-6} \otimes \mathbb{V} & \longrightarrow & 0.
 \end{array}$$

Remaining question: cohomology not clear. [no K operators, such that $S = DK - KD$.]

BGG in a more general setting (Čap, KH 2022)

$$\begin{array}{ccccccc}
 0 & \longrightarrow & Z^{0,0} & \xrightarrow{d^{0,0}} & Z^{1,0} & \xrightarrow{d^{1,0}} & \dots & \xrightarrow{d^{n-1,0}} & Z^{n,0} & \longrightarrow & 0 \\
 & & \uparrow K & \nearrow S^{0,1} & \uparrow K & \nearrow S^{1,1} & \uparrow K & \nearrow S^{n-1,1} & \uparrow K & & \\
 0 & \longrightarrow & Z^{0,1} & \xrightarrow{d^{0,1}} & Z^{1,1} & \xrightarrow{d^{1,1}} & \dots & \xrightarrow{d^{n-1,1}} & Z^{n,1} & \longrightarrow & 0 \\
 & & \uparrow K & \nearrow S^{0,2} & \uparrow K & \nearrow S^{1,2} & \uparrow K & \nearrow S^{n-1,2} & \uparrow K & & \\
 \dots & & \dots & & \dots & & \dots & & \dots & & \\
 & & \uparrow K & \nearrow S^{0,N} & \uparrow K & \nearrow S^{1,N} & \uparrow K & \nearrow S^{n-1,N} & \uparrow K & & \\
 0 & \longrightarrow & Z^{0,N} & \xrightarrow{d^{0,N}} & Z^{1,N} & \xrightarrow{d^{1,N}} & \dots & \xrightarrow{d^{n-1,N}} & Z^{n,N} & \longrightarrow & 0
 \end{array}$$

Assumptions:

$$\begin{aligned}
 S^{i+1,j-1} \circ S^{i,j} &= 0, \quad \forall i, j \geq 0, \\
 S^{i,j} &= d^{i,j-1} K^{i,j} - K^{i+1,j} d^{i,j} \quad S^{i,j-1} K^{i,j} = K^{i+1,j-1} S^{i,j}.
 \end{aligned}$$

Corollary: $dS = -Sd$.

- \rightarrow : de-Rham complex, \nearrow : (V^\bullet, S^\bullet) , Lie algebra (co)homology
- algebraic Hodge decomposition

$$Z^{i,j} = \mathcal{R}(S^{i-1,j+1}) \oplus \mathcal{R}(S^{i-1,j+1})^\perp = \mathcal{R}(S^{i-1,j+1}) \oplus \mathcal{N}(S^{i,j})^\perp \oplus H^{i,j},$$

where $H^{i,j} := \mathcal{R}(S^{i-1,j+1})^\perp \cap \mathcal{N}(S^{i,j})$ is an analogy of the space of harmonic forms.

These “algebraic harmonic forms” will be the spaces in the BGG sequences.

conformal deformation complex:

$$\begin{array}{ccccccc}
 0 & \longrightarrow & H^{q_0} \otimes V & \xrightarrow{\text{grad}} & H^{q_0-1} \otimes M & \xrightarrow{\text{curl}} & H^{q_0-2} \otimes M & \xrightarrow{\text{div}} & H^{q_0-3} \otimes V & \longrightarrow & 0 \\
 & & \nearrow S^{0,1} & & \nearrow S^{1,1} & & \nearrow S^{2,1} & & & & \\
 0 & \longrightarrow & H^{q_1} \otimes (\mathbb{R} \oplus V) & \xrightarrow{\text{grad}} & H^{q_1-1} \otimes (V \oplus M) & \xrightarrow{\text{curl}} & H^{q_1-2} \otimes (V \oplus M) & \xrightarrow{\text{div}} & H^{q_1-3} \otimes (\mathbb{R} \oplus V) & \longrightarrow & 0 \\
 & & \nearrow S^{0,2} & & \nearrow S^{1,2} & & \nearrow S^{2,2} & & & & \\
 0 & \longrightarrow & H^{q_2} \otimes V & \xrightarrow{\text{grad}} & H^{q_2-1} \otimes M & \xrightarrow{\text{curl}} & H^{q_2-2} \otimes M & \xrightarrow{\text{div}} & H^{q_2-3} \otimes V & \longrightarrow & 0.
 \end{array}$$

$$0 \longrightarrow H^q \otimes V \xrightarrow{\text{dev def}} H^{q-1} \otimes (\mathbb{S} \cap \mathbb{T}) \xrightarrow{\text{cot}} H^{q-4} \otimes (\mathbb{S} \cap \mathbb{T}) \xrightarrow{\text{div}} H^{q-5} \otimes V \longrightarrow 0$$

conformal Hessian complex:

$$\begin{array}{ccccccc}
 0 & \longrightarrow & H^{q_0} \otimes \mathbb{R} & \xrightarrow{\text{grad}} & H^{q_0-1} \otimes V & \xrightarrow{\text{curl}} & H^{q_0-2} \otimes V & \xrightarrow{\text{div}} & H^{q_0-3} \otimes \mathbb{R} & \longrightarrow & 0 \\
 & & \nearrow x \cdot & \nearrow I & \nearrow 2 \text{ vskw} & \nearrow x \cdot & \nearrow \text{tr} & \nearrow x \cdot & & & \\
 0 & \longrightarrow & H^{q_1} \otimes V & \xrightarrow{\text{grad}} & H^{q_1-1} \otimes M & \xrightarrow{\text{curl}} & H^{q_1-2} \otimes M & \xrightarrow{\text{div}} & H^{q_1-3} \otimes V & \longrightarrow & 0 \\
 & & \nearrow x \otimes & \nearrow I & \nearrow \text{rskw} & \nearrow x \otimes & \nearrow \frac{1}{3} \text{tr} & \nearrow x \otimes & & & \\
 0 & \longrightarrow & H^{q_2} \otimes \mathbb{R} & \xrightarrow{\text{grad}} & H^{q_2-1} \otimes V & \xrightarrow{\text{curl}} & H^{q_2-2} \otimes V & \xrightarrow{\text{div}} & H^{q_2-3} \otimes \mathbb{R} & \longrightarrow & 0.
 \end{array}$$

$$0 \longrightarrow H^q \otimes \mathbb{R} \xrightarrow{\text{dev hess}} H^{q-2} \otimes (\mathbb{S} \cap \mathbb{T}) \xrightarrow{\text{curl}} H^{q-3} \otimes (\mathbb{S} \cap \mathbb{T}) \xrightarrow{\text{div div}} H^{q-5} \otimes \mathbb{R} \longrightarrow 0$$

M: matrix

$$\begin{array}{ccccccc}
0 & \longrightarrow & H^q \otimes \mathbb{V} & \xrightarrow{\text{grad}} & H^{q-1} \otimes \mathbb{M} & \xrightarrow{\text{rot}} & H^{q-2} \otimes \mathbb{V} \longrightarrow 0 \\
& & & \searrow -\perp & & \nearrow \text{tr} & \\
0 & \longrightarrow & H^{q-1} \otimes (\mathbb{R} \times \mathbb{R}) & \xrightarrow{\text{grad}} & H^{q-2} \otimes (\mathbb{V} \times \mathbb{V}) & \xrightarrow{\text{rot}} & H^{q-3} \otimes (\mathbb{R} \times \mathbb{R}) \longrightarrow 0 \\
& & & \searrow -\iota & & \nearrow I & \\
0 & \longrightarrow & H^{q-2} \otimes \mathbb{V} & \xrightarrow{\text{grad}} & H^{q-3} \otimes \mathbb{M} & \xrightarrow{\text{rot}} & H^{q-4} \otimes \mathbb{V} \longrightarrow 0.
\end{array}$$

$$0 \longrightarrow H^q \otimes \mathbb{V} \xrightarrow{D^0} \begin{pmatrix} H^{q-1} \otimes (\mathbb{S} \cap \mathbb{T}) \\ H^{q-3} \otimes (\mathbb{S} \cap \mathbb{T}) \end{pmatrix} \xrightarrow{D^1} H^{q-3} \longrightarrow 0,$$

where, with pseudo-inverse T^\bullet ,

$$D^0 := \begin{pmatrix} \text{dev def} \\ \text{grad } T^{1,1} \text{ grad } T^{1,0} \text{ grad} \end{pmatrix}, \quad D^1 = \begin{pmatrix} \text{rot } T^{2,1} \text{ rot } T^{2,0} \text{ rot} \\ \text{rot} \end{pmatrix}^T.$$

Fixing 2D conformal Korn inequality:

$$\|u\|_3 \leq C(\|\text{dev def } u\|_2 + \|\text{grad } T^{1,1} \text{ grad } T^{1,0} \text{ grad } u\|), \quad \forall u \perp \mathcal{N}.$$

$$\text{dev def } u = \begin{pmatrix} \frac{1}{2}(\partial_x u_1 - \partial_y u_2) & \frac{1}{2}(\partial_y u_1 + \partial_x u_2) \\ \frac{1}{2}(\partial_y u_1 + \partial_x u_2) & -\frac{1}{2}(\partial_x u_1 - \partial_y u_2) \end{pmatrix}$$

Cauchy-Riemann operator. 2D conformal inequality with only C-R operator is invalid (Dain 2006). It can be fixed by including a third order term (Möbius structure).

Conclusion: cohomology of these complexes is isomorphic to the de Rham cohomology ✓

BGG is more than the complexes - it also provides machinery to derive results for elasticity, geometry... from results for de Rham.

Example: bounded Poincaré operators.

Poincaré operators: definition and motivation

Poincaré operators:

$$\cdots \rightleftarrows V^{i-1} \xrightleftharpoons[P^i]{d^{i-1}} V^i \xrightleftharpoons[P^{i+1}]{d^i} V^{i+1} \rightleftarrows \cdots$$

$P^k : V^k \mapsto V^{k-1}$, satisfying null-homotopy property

$$d^{k-1}P^k + P^{k+1}d^k = I_{V^k},$$

Motivation 1: constructing exact sequences and finite elements

(Hiptmair 1999, Arnold, Falk, Winther 2006)

$$du = 0 \implies u = (dP + Pd)u = d(Pu).$$

e.g., local exactness of de-Rham complexes.

Examples:

$$\cdots \longrightarrow \mathcal{P}_{r-(k-1)}\Lambda^{k-1} \xrightarrow{d^{k-1}} \mathcal{P}_{r-k}\Lambda^k \xrightarrow{d^k} \mathcal{P}_{r-(k+1)}\Lambda^{k+1} \longrightarrow \cdots$$

$$\cdots \longrightarrow \mathcal{P}_r\Lambda^{k-1} + P^k\mathcal{P}_r\Lambda^k \xrightarrow{d^{k-1}} \mathcal{P}_r\Lambda^k + P^{k+1}\mathcal{P}_r\Lambda^{k+1} \xrightarrow{d^k} \mathcal{P}_r\Lambda^k + P^{k+1}\mathcal{P}_r\Lambda^{k+1} \longrightarrow \cdots$$

Motivation 2: p -robustness of finite element methods

bounded, polynomial-preserving Poincaré operators imply results that are uniform with the polynomial degree.

Motivation 3: well-posedness of Stokes problem

given $f \in L^2$, find $u \in [H_0^1]^n$ and $p \in L^2/\mathbb{R}$, such that

$$\begin{aligned} -\Delta u + \nabla p &= f, \\ \operatorname{div} u &= 0. \end{aligned}$$

inf-sup condition: for any $q \in L^2/\mathbb{R}$, $\exists u = P(q) \in [H_0^1]^n$, s.t. $\operatorname{div} u = q$, $\|u\|_1 \leq C\|q\|$.

Motivation 4: *analytic results*, e.g., regular decomposition, compactness.

How to construct

Smooth de-Rham complex (see books on manifolds)

Let $F_t : \Omega \rightarrow \Omega$, $t \in [0, 1]$ be a continuous family of operators indexed by t .

If u is a k -form:

$$(\mathfrak{p}[F]u)_x(\xi_2, \dots, \xi_k) = \int_0^1 u_{F_t(x)}(\partial_t F_t(x), DF_t(x)\xi_2, \dots, DF_t(x)\xi_k) dt.$$

Suppose that F_1 is identity and F_0 is constant x_0 , then we have $d\mathfrak{p} + \mathfrak{p}d = I$ for $k \geq 1$ and $\mathfrak{p}du(x) = u(x) - u(x_0)$ for $k = 0$.

Simplification

- 1D with *base point* x_0

$$\mathfrak{p}(u) := \int_{x_0}^x u(y) dy, \quad \partial \mathfrak{p}(u) = u, \quad \mathfrak{p}(\partial v) = v(x) - v(x_0).$$

- 3D vector proxy (choose a curve $\gamma(t) = tx$ connecting 0 and x):

$$\mathfrak{p}^1 u = \int_0^1 u_{tx} \cdot x dt, \quad \mathfrak{p}^2 v = \int_0^1 tv_{tx} \wedge x dt, \quad \mathfrak{p}^3 w = \int_0^1 t^2 w_{tx} x dt.$$

Sobolev de-Rham complex: Costabel-McIntosh 2010

$$0 \longrightarrow H^q \Lambda^0 \xrightarrow{d^0} H^{q-1} \Lambda^1 \xrightarrow{d^1} \dots \xrightarrow{d^{n-1}} H^{q-n} \Lambda^n \longrightarrow 0,$$

where q is any real number.

- regularized Poincaré operators: averaging the base points in the smooth de-Rham version, mapping $H^q \Lambda^k$ to $H^{q+1} \Lambda^{k-1}$, polynomial-preserving
- generalized Bogolovskiĭ operators: “dual” of Poincaré, , mapping $H_0^q \Lambda^k$ to $H_0^{q+1} \Lambda^{k-1}$

pseudo-differential operators of order -1, which implies boundedness between various spaces (Sobolev, Besov...)

General Lipschitz domain:

$$dP + Pd = I - L,$$

where $L(u) \in C^\infty$ for any u . The smoothing operator L comes from partition of unity (into a union of star-shaped patches).

Goal: BGG version

Outline:

$$\begin{array}{ccc}
 Y^i & \xrightleftharpoons[d^i]{P^{i+1}} & Y^{i+1} \\
 F \downarrow & & F \downarrow \\
 Y^i & \xrightleftharpoons[d_V^i]{P_V^{i+1}} & Y^{i+1} \\
 B^i \updownarrow A^i & & B^{i+1} \updownarrow A^{i+1} \\
 \Upsilon^i & \xrightleftharpoons[\mathcal{D}^i]{\mathcal{D}^{i+1}} & \Upsilon^{i+1}
 \end{array}$$

(Y^\bullet, d^\bullet) : de Rham complex
 (Y^\bullet, d_V^\bullet) : twisted complex
 $(\Upsilon^\bullet, \mathcal{D}^\bullet)$: BGG complex

Poincaré operators for BGG complexes:

$$\mathcal{P} := B \circ F \circ P \circ F^{-1} \circ A.$$

If P^\bullet satisfies $dP + Pd = I$, then \mathcal{P}^\bullet satisfies $\mathcal{D}\mathcal{P} + \mathcal{P}\mathcal{D} = I$.

$$F^i := (I - P^{i+1}S^i)^{-1} = \sum_{l=0}^{\infty} (P^{i+1}S^i)^l$$

$$\begin{array}{ccccccc}
0 & \longrightarrow & Z^{0,0} & \begin{array}{c} \xleftarrow{d^{0,0}} \\ \xrightarrow{S^{0,1}P^{1,0}} \end{array} & Z^{1,0} & \begin{array}{c} \xleftarrow{d^{1,0}} \\ \xrightarrow{S^{1,1}P^{2,0}} \end{array} & \dots & \begin{array}{c} \xleftarrow{d^{n-1,0}} \\ \xrightarrow{S^{n-1,1}P^{n,0}} \end{array} & Z^{n,0} & \longrightarrow & 0 \\
0 & \longrightarrow & Z^{0,1} & \begin{array}{c} \xleftarrow{d^{0,1}} \\ \xrightarrow{S^{0,2}P^{1,1}} \end{array} & Z^{1,1} & \begin{array}{c} \xleftarrow{d^{1,1}} \\ \xrightarrow{S^{1,2}P^{2,1}} \end{array} & \dots & \begin{array}{c} \xleftarrow{d^{n-1,1}} \\ \xrightarrow{S^{n-1,2}P^{n,1}} \end{array} & Z^{n,1} & \longrightarrow & 0 \\
\dots & & \dots & & \dots & & & & \dots & & \\
0 & \longrightarrow & Z^{0,N} & \begin{array}{c} \xleftarrow{d^{0,N}} \\ \xrightarrow{S^{0,N}P^{1,N}} \end{array} & Z^{1,N} & \begin{array}{c} \xleftarrow{d^{1,N}} \\ \xrightarrow{S^{1,N}P^{2,N}} \end{array} & \dots & \begin{array}{c} \xleftarrow{d^{n-1,N}} \\ \xrightarrow{S^{n-1,N}P^{n,N}} \end{array} & Z^{n,N} & \longrightarrow & 0.
\end{array}$$

Twisted complexes

$$\dots \longrightarrow Y^{i-1} \xrightarrow{d_V^{i-1}} Y^i \xrightarrow{d_V^i} Y^{i+1} \longrightarrow \dots,$$

$$Y^i := \begin{pmatrix} Y^{i,0} \\ Y^{i,1} \\ \vdots \\ Y^{n-1,0} \\ Y^{n,0} \end{pmatrix}, \quad d_V := d - S = \begin{pmatrix} d^{i,0} & -S^{i,1} & & & \\ & d^{i,1} & -S^{i,2} & & \\ & & \ddots & \ddots & \\ & & & d^{i,n-1} & -S^{i,n} \\ & & & & d^{i,n} \end{pmatrix}$$

- twisted complexes: Timoshenko beam, Reissner-Mindlin plate, Cosserat elasticity
- BGG complexes: Euler-Bernoulli beam, Kirchhoff-Love plate, standard elasticity 14 / 19

1D example

BGG diagram

$$\begin{array}{ccc} H^q & \xrightarrow{\partial} & H^{q-1} \\ & \nearrow I & \\ H^{q-1} & \xrightarrow{\partial} & H^{q-2}, \end{array}$$

where $\partial := \frac{d}{dx}$.

BGG complex:

$$H^q \xrightarrow{\partial^2} H^{q-2}$$

Poincaré operators: with $P_{\sharp} : H^{q-1} \rightarrow H^q$ and $P_b : H^{q-2} \rightarrow H^{q-1}$.

$$\mathcal{P} := P_{\sharp} P_b.$$

satisfying

$$\partial^2 \mathcal{P} = I_{H^{q-2}}, \quad \mathcal{P} \partial^2 = I_{H^q} + \alpha,$$

where $\alpha \in \mathcal{N}(\partial^2)$.

3D elasticity complex

$\mathcal{P}^1 = P(PS - SP)Td$ (all ingredients given in BGG diagram).

For *smooth* e (Christiansen, KH, Sande 2019):

$$\mathcal{P}^1(e) = \int_0^x e(y) \cdot dy - \underbrace{\int_0^x (\text{mskw} \int_0^y \underbrace{(e(z) \times \nabla)^T}_{Td(e)} \cdot dz) \cdot dy}_{SPTd(e)}.$$

For *general Sobolev functions*: generalization of Cesàro-Volterra formula (1906, 1907), satisfying

$$\mathcal{P}^1 \text{def}(w) = w(x) - w(0) + \frac{1}{2} \int_0^x dy \wedge \nabla w(y).$$

Compared to

- 'A Cesàro-Volterra formula with little regularity', Ciarlet, Gratie, Mardare, 2010 JMPA,

the new formulas are explicit, polynomial-preserving, work for a broad class of functions, for the entire complex.

Complex property $P \circ P = 0$?

Let P^\bullet be Poincaré operators satisfying $dP + Pd = I$, but not necessarily $P \circ P = 0$. We can generally modify P^\bullet to \tilde{P}^\bullet , defined by $\tilde{P} := P - DP^2 = PDP$. Then straightforward algebra implies $d\tilde{P} + \tilde{P}d = I$ and $\tilde{P} \circ \tilde{P} = 0$.

Nontrivial cohomology?

Standard procedure: cover $\Omega = \cup_j \Omega_j$, partition of unity, $1 = \sum_j \xi_j$ with ξ_j supported on Ω_j . Use $dP + Pd = I$ on each Ω_j , leading to $dP + Pd = I - L$ globally, where L is a smoothing operator.

Polynomial-preservation?

BGG Poincaré operators are polynomial-preserving if the de Rham versions are so (e.g., regularized Poincaré).

Some applications

Construction of polynomial exact sequences, e.g.,

$$\text{RM} \longrightarrow \mathcal{P}_r \otimes \mathbb{V} \xrightarrow{\text{def}} \mathcal{P}_{r-1} \otimes \mathbb{S} \xrightarrow{\text{inc}} \mathcal{P}_{r-3} \otimes \mathbb{S} \xrightarrow{\text{div}} \mathcal{P}_{r-4} \otimes \mathbb{V} \longrightarrow 0.$$

One may also derive \mathcal{P}_r^- type spaces (Raviart-Thomas etc.), which have been applied to constructing numerical schemes (Zhao et al. 2023), e.g.,

$$0 \longrightarrow \mathcal{P}_r \otimes \mathbb{V} + \mathcal{P}^1[\mathcal{P}_r \otimes \mathbb{S}] \xrightarrow{\text{def}} \mathcal{P}_r \otimes \mathbb{S} + \mathcal{P}^2[\mathcal{P}_r \otimes \mathbb{S}] \xrightarrow{\text{inc}} \mathcal{P}_r \otimes \mathbb{S} + \mathcal{P}^3[\mathcal{P}_r \otimes \mathbb{V}] \xrightarrow{\text{div}} \mathcal{P}_r \otimes \mathbb{V} \longrightarrow 0.$$

hp finite elements.

For any $v \in L^2 \otimes \mathbb{V}$, $\exists \sigma = \mathcal{P}(v) \in H^1(\mathbb{S})$, s.t., $\text{div } \sigma = v$, and \mathcal{P} is polynomial-preserving.

\implies *hp* stability for FEM for Hellinger-Reissner formulation (Aznanan, KH, Parker, in preparation)

References:

- *Complexes from complexes*, Douglas Arnold, KH; *Foundations of Computational Mathematics* (2021). [complexes from two rows](#)
- *BGG sequences with weak regularity and applications*, Andreas Čap, KH; *to appear*, *Foundations of Computational Mathematics* (2022) [multi-rows](#)
- *Poincaré path integrals for elasticity*, Snorre Christiansen, KH, Espen Sande, *Journal de Mathématiques Pures et Appliquées*, (2019) [smooth elasticity complex in 3D](#)
- *Bounded Poincaré operators for BGG complexes*, Andreas Čap, KH; *to appear*, *Journal de Mathématiques Pures et Appliquées* (2023) [general BGG complexes, Sobolev spaces](#)